

Fluid mechanics

Q. Sec. A. 1 (i) c (ii) c (iii) c (iv) c (v) c (vi) a (vii) a  
(viii) a (ix) c (x) c.

Sec. B Q. 2 (i) Eddy Viscosity :- eddy viscosity  $E_v$  is analogous to  $\mu$ , the absolute viscosity. The relation ship bet<sup>n</sup> shear stress and velocity gradient in a turbulent stream is used to define an eddy viscosity.

$$\tau_{tge} = E_v \frac{du}{dy}$$

(ii) Boundary layer :- Boundary layer is defined as that part of a moving fluid in which the fluid motion is influenced by the presence of a solid boundary.  
Short explanation of boundary layer formation may be given.

(iii) Transition length :- The length of the entrance region of the tube necessary for the boundary layer to reach the center of the tube and for fully developed flow to be established is called the transition length.

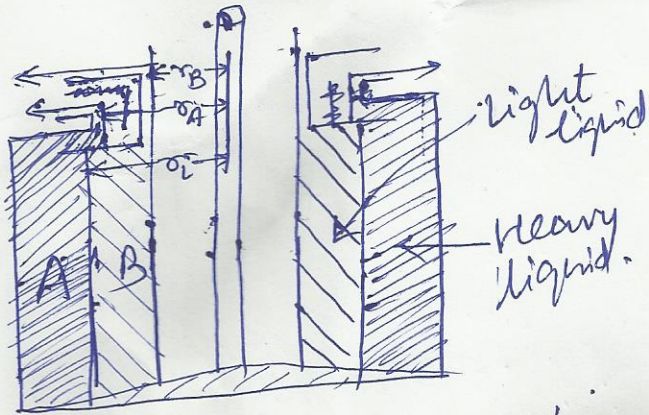
$$\frac{x_t}{D} \approx 0.05 N_{Re} \quad (\text{for laminar flow})$$

(iv) Hydrostatic law :- In a stationary mass of a single static fluid, the pressure is constant in any cross section  $\parallel$  to the earth's surface but varies from height to height.

$P = \rho g z$  is the mathematical expression (2)

or

(a) Centrifugal decanter :- It is used when the density difference of the two liquids is so small.



$$P_i - P_B = P_i - P_A$$

$$P_i - P_B = \frac{\omega^2 \rho_B (r_i^2 - r_B^2)}{2g_c}$$

$$P_i - P_A = \frac{\omega^2 \rho_A (r_i^2 - r_A^2)}{2g_c}$$

$$\therefore \rho_B (r_i^2 - r_B^2) = \rho_A (r_i^2 - r_A^2)$$

$$\therefore r_i = \sqrt{\frac{r_A^2 - (\rho_B/\rho_A) r_B^2}{1 - \rho_B/\rho_A}}$$

(b)  $\omega = \pi \times \frac{4000}{60} = 209.4 \text{ rad/s.}$

$\rho = 1109 \text{ kg/m}^3.$

$r_2 = 0.125 \text{ m}; \quad r_1 = 0.125 - 0.05 = 0.075.$

$P_2 - P_1 = \frac{\omega^2 \rho (r_2^2 - r_1^2)}{2g_c}$

$= \frac{(209.4)^2 \times 1109 \times [(0.125)^2 - (0.075)^2]}{2}$

$= 243139$

$\text{N/m}^2.$

$0.125 - 0.05 = 0.075$

$0.010000$

Q.3.  $A_1 v_1 = A_2 v_2$   $v_2 = 20 \text{ m/s}$ ,  $Q = 10 \text{ m}^3/\text{s}$ . (3)

$$\frac{P_1}{\rho g} + \frac{v_1^2}{2g} = \frac{P_2}{\rho g} + \frac{v_2^2}{2g}$$

$$\frac{2.943 \times 10^4}{1.16 \times 9.81} + \frac{(10)^2}{2 \times 9.81} = \frac{P_2}{\rho g} + \frac{(20)^2}{2 \times 9.81}$$

$$P_2 = 29.25 \text{ kN/m}^2$$

$$F_x = \rho Q [v_{1x} - v_{2x}] + (P_1 A_1)_x + (P_2 A_2)_x$$

$$(P_1 A_1)_x = P_1 A_1 = 29.43 \text{ kN/m}^2 \times 1 = 29.43 \text{ kN}$$

$$(P_2 A_2)_x = P_2 A_2 \cos 45 = -29.25 \frac{\text{kN}}{\text{m}^2} \times 0.5 \times 0.7071$$

$$F_x = +19038.59 \text{ N}$$

$$F_y = \rho Q [v_{1y} - v_{2y}] + (P_1 A_1)_y + (P_2 A_2)_y$$

$$v_{1y} = 0, v_{2y} = v_2 \sin 45 = 20 \times 0.7071 = 14.142$$

$$(P_1 A_1)_y = 0, (P_2 A_2)_y = -P_2 A_2 \sin 45 = -29255.8 \times 0.5 \times \sin 45 = -10343.37$$

$$F_y = -10507.42 \text{ N}$$

$$\text{Resultant force } F_R = \sqrt{F_x^2 + F_y^2} = 21746.6 \text{ N}$$

OR.

Bernoulli's theorem :- Euler's eqn<sup>n</sup> of motion -  
consider a stream line in which flow is taking place in s-direction. Consider a

Cylindrical element of C.S.  $dA$  and length  $ds$  @

The force balance -

$$PdA - \left(P + \frac{\partial P}{\partial s} ds\right)dA - \rho g dA ds \cos \theta.$$

$$= \rho dA ds a_s.$$

$a_s$  - accelera<sup>n</sup> in  $s$  direc<sup>n</sup>.

$$a_s = \frac{dv}{dt} = \frac{\partial v}{\partial s} \frac{ds}{dt} + \frac{\partial v}{\partial t} = v \frac{\partial v}{\partial s} + \frac{\partial v}{\partial t}$$

If the flow is steady,  $\frac{\partial v}{\partial t} = 0$ .

$$a_s = v \frac{dv}{ds}$$

$$\therefore \frac{\partial P}{\partial s} ds dA - \rho g dA ds \cos \theta = \rho dA ds v \frac{dv}{ds}$$

Dividing by  $\rho ds dA$ ;  $-\frac{\partial P}{\rho \partial s} - g \cos \theta = v \frac{dv}{ds}$

$$\text{Or } \frac{\partial P}{\rho \partial s} + g \cos \theta + v \frac{dv}{ds} = 0.$$

But  $\cos \theta = \frac{dz}{ds}$ .

$$\therefore \frac{1}{\rho} \frac{\partial P}{\partial s} + g \frac{dz}{ds} + \frac{v dv}{ds} = 0 \quad \text{or } \frac{\partial P}{\rho} + g dz + v dv = 0$$

$$\text{Or } \frac{\partial P}{\rho} + g dz + v dv = 0$$

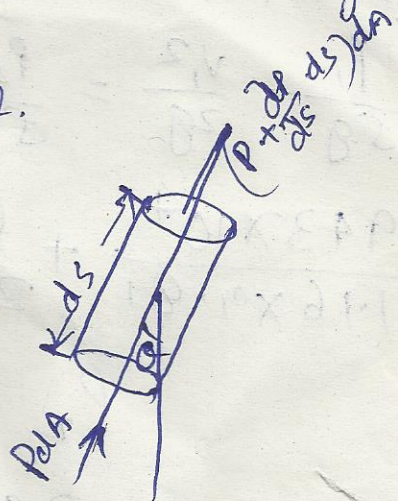
Bernoulli's eqn<sup>n</sup> is obtained by integrating the Euler's eqn<sup>n</sup> of motion as

$$\int \frac{\partial P}{\rho} + \int g dz + \int v dv = \text{Const.}$$

If flow is incompressible,  $\rho$  is const. and,

$$\frac{P}{\rho} + g z + \frac{v^2}{2} = \text{Const.}$$

$$\text{Or } \frac{P}{\rho g} + z + \frac{v^2}{2g} = C \quad \text{Bernoulli's eqn<sup>n</sup>}$$



Q. 4. Hagen Poiseuille eqn. :-  $\Delta P_s = \frac{32 \Delta L \bar{V} \mu}{g_c D^2}$  (5)

For laminar flow of newtonian fluids in circular pipe, the local velo.  $u$  depends only on the radius  $r$ . Also, the element of area  $ds$  is that of a thin ring of radius  $r$  with width  $dr$ . The area of this elementary ring is

$$ds = 2\pi r dr.$$

The desired velocity distribution relation is a function of  $r$ .

A direct method of obtaining the velo. distribution for newtonian fluids is to use the definition of viscosity.

$$\mu = - \frac{\tau_{rc}}{du/dr} \quad \text{--- (1)}$$

as we know,  $\frac{\tau_{rc}}{r\omega} = \frac{\tau}{r}$  --- (2)

$$\therefore \frac{du}{dr} = - \frac{\tau_{rc}}{\mu} = - \frac{\tau_{rc}}{r\omega\mu} \cdot r \quad \text{--- (3)}$$

Integrat<sup>n</sup> of (3) with the B.C.  $u=0$ ,  $r=r_w$ ; gives

$$\int_0^u du = - \frac{\tau_w g_c}{r_w \mu} \int_{r_w}^r r dr$$

$$u = \frac{\tau_w g_c}{2 r_w \mu} (r_w^2 - r^2)$$

$$u_{max} = \frac{\tau_w g_c r_w}{2 \mu} \quad [r=0]$$

Then  $\tau_w$  for average velo.  $\bar{V}$ ,

we use,  $ds = 2r \, dr$  (6)

$$u = \frac{\tau_w g_c}{2 \rho_w \mu} (\rho_w^2 - r^2)$$

$$S = \pi \rho_w^2$$

$$\bar{V} = \frac{1}{S} \int_S u \, ds$$

$$= \frac{1}{\pi \rho_w^2} \int_S \frac{\tau_w g_c}{2 \rho_w \mu} (\rho_w^2 - r^2) \cdot 2r \, dr$$

$$= \frac{\tau_w g_c}{\rho_w^3 \mu} \int_0^{\rho_w} (\rho_w^2 - r^2) r \, dr$$

$$\bar{V} = \frac{\tau_w g_c \rho_w}{4 \mu} \quad \text{--- (4)}$$

Now, we transform this eqn. (4) by eliminating  $\tau_w$  in favour of  $\Delta P_S$  by using eqn.

$$\tau_w = \frac{2 \rho_w \Delta L}{8 \rho_w} = \frac{\Delta P_S}{8} \Rightarrow 4f \frac{\Delta L \bar{V}^2}{2D g_c}$$

$$\bar{V} = \frac{\Delta P_S g_c \rho_w}{\Delta L} \frac{\rho_w}{2} \frac{\rho_w}{4 \mu} = \frac{\Delta P_S g_c D^2}{32 \Delta L \mu}$$

$$\therefore \Delta P_S = \frac{32 \Delta L \bar{V} \mu}{g_c D^2}$$

Limitations :- applicable only for laminar flow, Newtonian fluid, circular channel.

Significance - Used for practical calculation of viscosity.

(a)

(7)

$$\mu = \frac{\tau_{ge}}{du/dr}, \quad \text{or} \quad \frac{\tau_w}{r_w} = \frac{\tau}{r}$$

$$\frac{du}{dr} = -\frac{\tau_{ge}}{r} = -\frac{\tau_w r_{ge}}{r_w \mu} \cdot r$$

$$\int_0^h du = -\frac{\tau_w r_{ge}}{r_w \mu} \int_{r_w}^r r dr$$

$$u = -\frac{\tau_w r_{ge}}{2 r_w \mu} (r_w^2 - r^2)$$

$$u_{max} = \frac{\tau_w r_{ge} r_w}{2 \mu} \quad [r=0]$$

for average vel.

$$ds = 2\pi r dr, \quad u = \frac{\tau_w r_{ge}}{2 r_w \mu} (r_w^2 - r^2)$$

$$s = \pi r_w^2, \quad \bar{V} = \frac{1}{s} \int u ds$$

$$\bar{V} = \frac{1}{\pi r_w^2} \int_0^{r_w} \frac{\tau_w r_{ge}}{2 r_w \mu} (r_w^2 - r^2) 2\pi r dr$$

$$\bar{V} = \frac{2}{3} u_{max}$$

$$\frac{u}{u_{max}} = 1 - (r/r_w)^2$$

(b) Fanning friction factor - which is a common parameter especially useful in the study of turbulent flow. It is denoted by  $f$ , and may be defined as the ratio of the wall shear stress to the product of the density and the velo. head

$$f = \frac{\tau_w}{\rho V^2 / 2g} = \frac{2g \tau_w}{\rho V^2}$$

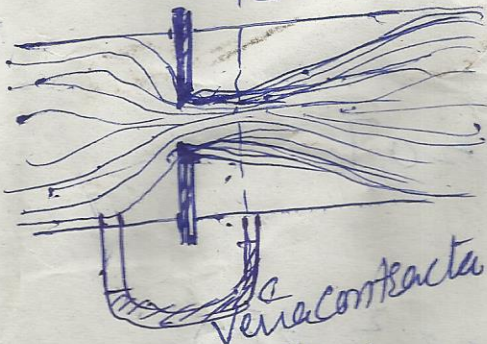
$$\frac{h_f}{\rho g} = \frac{2Lw}{\rho g v} \Delta L = \frac{\Delta P_f}{\rho} = 24 f \frac{\Delta L v^2}{D \rho g}$$

$$f = \frac{\Delta P_f \rho g D}{2 \Delta L \rho v^2}$$

(5)

### Unit IV

(5) (a) Orifice meter -



It is a variable head flow meter. It meets the objections to the venturi but at the price of a larger power consumption. It consists of an accurately

machined & drilled plate mounted bet<sup>n</sup>. two flanges with the hole concentric with the pipe in which it is mounted. Press. taps, one above and one below the orifice plate are installed and connected to a manometer. The reduction of the cross section of the flowing stream in passing through the orifice increases the velocity head at the expense of the press. head and the reduction in pressure between the taps is measured by the manometer. Bernoulli's eqn<sup>n</sup> provides a basis for correlating the increase in velocity head with the decrease in press. head.

Because of the sharpness of the orifice, the fluid stream separates from the downstream side of the orifice plate and forms a free flowing jet in the downstream fluid side fluid; a vena contracta forms. The jet is not under the



contraction of solid walls, and the area of the (a) jet varies from that of the opening to that of the vena contracta. The area at the vena-contracta is minimum.

$$u_0 = \frac{C_0}{\sqrt{1-\beta^4}} \sqrt{\frac{2g_c(h_a - h_b)}{5}}$$

(b) The press. drop in the upstream cone is utilized to measure the rate of flow and the decreased velocity is largely recovered in the downstream cone. To make the press. recovery large, the angle of the downstream cone is small so boundary layer separation is prevented and friction minimized. Since separation does not occur in a contracting cross section, the upstream cone can be made shorter than the downstream cone.

or

(i) NPSH and Cavitation:- If the suction press. is less <sup>or slightly higher</sup> than the vapor press., some liquid may flash to vapor inside the pump, this is called cavitation.

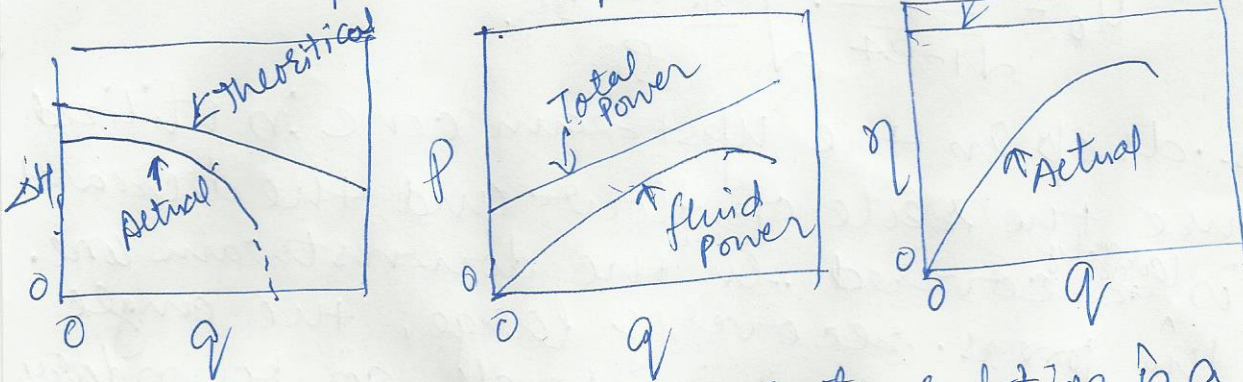
To avoid cavitation, the pressure at the pump inlet must exceed the vapor press. by a certain value, called the net positive suction head (NPSH). The required value of NPSH is about 2 to 3 m for small centrifugal pumps.

$$NPSH = \frac{f_c}{g} \left( \frac{P_a - P_v}{\rho} - h_{fs} \right) - Z_a$$

$P_a$  = absolute press. at surface of reservoir

$P_v$  = vap. press.,  $h_{fs}$  = friction in suction line.

(ii) Characteristics curves - The plots of actual head, total power consumption, and efficiency vs vol. flow rate are called the characteristics curves of a pump.



The theoretical head-flow rate relation is a straight line, the actual developed head is considerably less and drops to zero as the rate increases to a certain value in a given pump.

The difference between typical curves of fluid power and total power vs flow rate are shown. The difference between ideal & actual performance represents the power lost in the pump it results from fluid friction and shock losses.

The pump efficiency is the ratio of fluid power to the total power input. The curve shows that the efficiency rises rapidly with flow rates at low rates, reaches a maximum in the region of the rated capacity, then falls as the flow rate approaches the zero-head value.

~~Unit I~~      Unit II  
OR

6 (a)  $P = \frac{N p v^3 D_a^5}{g_c}$        $N_p = 5.8, \quad n = \frac{90}{60} = 1.5 \text{ r/s}$

$P = \frac{5.8 \times (1.5)^3 \times (1.49)^5}{32.2} = 2476.6 \text{ W}$

Power,  $\frac{1921}{580} \times 2.33 \text{ kW} = 2.47 \text{ kW}$       or       $5.8 \times (1.5)^3 \times 1.49^5 = 2.476 \text{ kW}$